

1 Wireless Communication with Nanoplasmonic Data Carriers: 2 Macroscale Propagation of Nanophotonic Plasmon Polaritons 3 Probed by Near-Field Nanoimaging

4 Moshik Cohen,^{*,†,§} Yossi Abulafia,[§] Dmitry Lev,^{||} Aaron Lewis,^{||} Reuven Shavit,[‡] and Zeev Zalevsky^{†,§}

5 [†]Faculty of Engineering, Bar-Ilan University, Ramat-Gan 52900, Israel

6 [‡]Department of Electrical and Computer Engineering, Ben-Gurion University of the Negev, Beer-Sheva 84105, Israel

7 [§]Bar-Ilan Institute for Nanotechnology & Advanced Materials, Ramat-Gan 52900, Israel

8 ^{||}Department of Applied Physics, Selim and Rachel Benin School of Engineering and Computer Science, The Hebrew University,
9 Givat Ram, Jerusalem 9190401, Israel

10 **ABSTRACT:** The ability to control the energy flow of light at the nanoscale is
11 fundamental to modern communication and big-data technologies, as well as quantum
12 information processing schemes. However, because photons are diffraction-limited, efforts
13 of confining them to dimensions of integrated electronics have so far proven elusive. A
14 promising way to facilitate nanoscale manipulation of light is through plasmon
15 polaritons—coupled excitations of photons and charge carriers. These tightly confined
16 hybrid waves can facilitate compression of optical functionalities to the nanoscale but suffer
17 from huge propagation losses that limit their use to mostly subwavelength scale
18 applications. With only weak evidence of macroscale plasmon polaritons, propagation has
19 recently been reported theoretically and indirectly, no experiments so far have directly
20 resolved long-range propagating optical plasmon polaritons in real space. Here, we launch
21 and detect nanoscale optical signals, for record distances in a wireless link based on novel
22 plasmonic nanotransceivers. We use a combination of scanning probe microscopies to
23 provide high resolution real space images of the optical near fields and investigate the long-
24 range propagation of nanoscale optical signals. We design our nanotransceivers based on a high-performance nanoantenna,
25 *Plantenna*, hybridized with channel plasmon waveguides with a cross-section of 20 nm × 20 nm, and observe propagation for
26 distances up to 1000 times greater than the plasmon wavelength. We experimentally show that our approach hugely outperforms
27 both waveguide and wireless nanophotonic links. This successful alliance between *Plantenna* and channel plasmon waveguides
28 paves the way for new generations of optical interconnects and expedites long-range interaction between quantum emitters and
29 photomolecular devices.



30 **KEYWORDS:** Plasmonics, nanoantennas, channel waveguides, wireless, nanoimaging

31 **T**he proposed scheme is designed to enable macroscale
32 communication between nanoscale devices utilizing
33 surface plasmon polaritons (SPPs). Hence, we use channel
34 waveguides that confine SPPs to their channel dimensions,
35 which can be as small as several nanometers.^{1–8} However, as
36 dimensions decrease, SPPs exhibit increased losses that limit
37 their propagation in waveguides to distances of only few
38 micrometers. To address this fundamental limitation, we
39 convert channel SPPs to optical surface waves that propagate
40 for significantly larger distances on dielectric substrates. A high-
41 efficiency nanoreceiver, designed to convert surface waves to
42 channel SPPs, is placed the remote edge of the system. **Figure**
43 **1a** illustrates the proposed communication nanosystem, which
44 (a) converts light to nanoscale SPPs, (b) propagates SPPs in
45 channel waveguide, (c) converts these SPPs to surface waves
46 and propagate them for long distance, and (d) excites SPPs
47 from the surface waves at remote locations. As shown in the
48 right-hand side of **Figure 1a**, laser light illuminates the
49 *Plantenna* to launch SPPs at the waveguide. Second, *Plantenna*,

located at the other edge of the waveguide, converts these SPPs 50
to surface waves that propagate on the substrate. The surface 51
waves are reconverted to SPPs at a remote, *Plantenna* based 52
nanoreceiver. We use waveguides with a propagation loss of 53
 $e^{-\alpha l}$, where the absorption constant $\alpha = (18 \mu\text{m})^{-1}$ for a 54
channel width of 20 nm at a red wavelength of $\lambda = 633 \text{ nm}$ and 55
 l is the propagation length. 56

In contrast, absorption for wireless links occur only at the 57
antennas and are much lower than for a waveguide. For 58
conventional (e.g., dipole, bowtie) nanoantennas, the prop- 59
agation loss for wireless links behaves like $(D/l)^2$, where D is 60
the directivity.⁹ Here, we show that *Plantenna* based wireless 61
links hugely outperform both waveguide and conventional 62
nanoantenna based alternatives. **Figure 1b** presents a 3D model 63
of the nanotransceiver, with zoom in to the *Plantenna* region 64

Received: January 19, 2017

Revised: April 25, 2017

Published: May 3, 2017

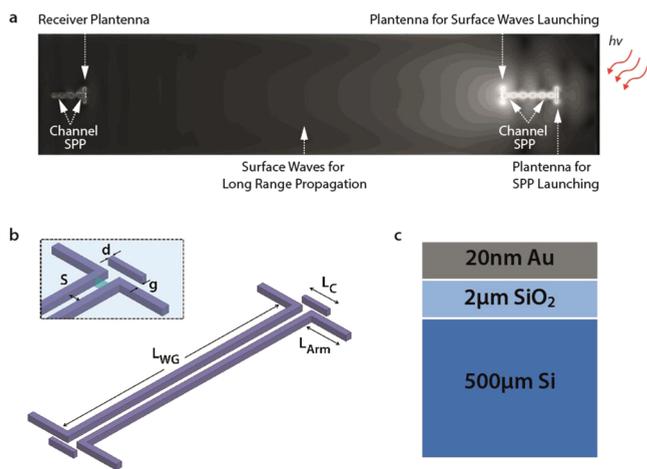


Figure 1. Wireless communications with optical plasmon polaritons. (a) Illustration of the proposed nanoscale communication system. SPP launching Plantenna (right) converts free space light to propagating waveguides SPPs, which are coupled to surface waves by the “Surface Waves Launching Plantenna” for long-range propagation. A Plantenna based nanoreceiver (left) converts the surface waves to channel SPPs at remote distances. (b) 3D model of a Plantenna based plasmonic nanotransceiver. Zoom in to the Plantenna region is shown in the inset c, materials stack up used to fabricate the devices.

65 shown in the inset. The physical principle behind the Plantenna
 66 invention is the enormous field enhancement and confinement
 67 exhibited by resonant, optically illuminated adjacent metallic
 68 nanoparticles. These properties, mainly originated from
 69 coherent capacitive coupling between the particles, are
 70 significantly better than those of isolated nanoparticles. The
 71 Plantenna comprised of two metallic nanorods of length L_{Arm} ,
 72 separated by a nanoscopic gap ($s = 10\text{--}35\text{ nm}$), in a dipole
 73 arrangement. An additional nanorod, termed director, is placed
 74 at much closer proximity of only 7 nm ($g \sim 7\text{ nm}$) to the
 75 dipole. A detailed analysis on the Plantenna physics, which also
 76 includes optimization for high efficiency excitation of channel
 77 SPPs, was recently reported.⁵ Figure 1c shows the material
 78 stack up used in this work, comprised of 20 nm Au layer
 79 deposited on a Si on insulator (SOI) wafer ($500\text{ }\mu\text{m Si}$, $2\text{ }\mu\text{m}$
 80 SiO_2), for potential CMOS computability.

81 For nanofabrication, we use electron beam lithography
 82 (EBL), ion beam sputtering (Au, 20 nm), and liftoff. After
 83 liftoff, the resist is completely removed, allowing contact mode
 84 near-field optical characterization. We fabricated devices
 85 comprised of standalone nanotransceivers and complete
 86 communication systems. Figure 2a shows a high-resolution
 87 scanning electron microscopy (HR-SEM) image of a fabricated
 88 nanotransceiver, recorded at beam current of 0.4 nA and low
 89 accelerating voltage of 5 kV, for sub 1 nm imaging resolution;
 90 corresponding 3D AFM topography is shown in Figure 2b.
 91 Nanotransceivers with dimensions of $L_{Arm} = 220\text{ nm}$, $L_C = 120$
 92 nm , $s = 20\text{ nm}$, $g = 7\text{ nm}$, and $L_{WG} = 1.5\text{ }\mu\text{m}$ were fabricated
 93 successfully and repeatedly. Figure 2c shows near-field KPFM
 94 nanoimaging under illumination with a He–Ne laser ($\lambda_0 = 633$
 95 nm), recorded at a set lift height of 30 nm using a high aspect
 96 ratio uncoated Si AFM tip with a diameter of 2 nm. As
 97 observed, the laser light is efficiently converted to propagating
 98 plasmons at the waveguide channel by the Rx (right) Plantenna
 99 and then recoupled to surface waves via the Tx (left) Plantenna.
 100 Characterized by periodic peaks (purple) in the KPFM signal
 101 imaged at the waveguide channel, SPPs with an effective

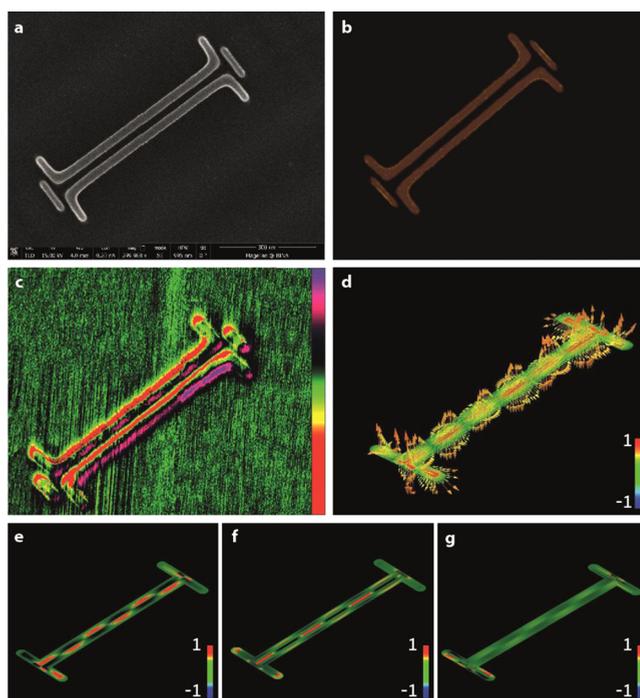


Figure 2. Plantenna-based plasmonic nanotransceiver. (a) High-resolution SEM image of the fabricated nanotransceiver. (b) 3D AFM image of the fabricated nanotransceiver. (c) KPFM under optical illumination analysis of the nanotransceiver. KPFM signal scale bar: $\pm 4.7\text{ V}$. (d) Numerically calculated optical near-field vector. (e) Numerically calculated optical near-field image showing $\text{Re}(E_z) = |E_z| \cos(\phi_z)$. (f) Numerically calculated optical near-field showing $\text{Re}(E_x)$. (g) Numerically calculated optical near field showing $\text{Re}(E_y)$. Scale bar: 100 nm.

wavelength of 35–150 nm were measured. The experimental
 102 results are reproduced by numerical calculation results,
 103 presented at the optical frequency of 474 THz (633 nm).
 104 The theoretical results are obtained using a high-frequency
 105 structure simulator based on the finite element method
 106 (FEM).^{3–5,10,11} Numerical calculation results of the device
 107 are shown in Figure 2d–g, with Figure 2d showing the local
 108 near-field vector in 3D, and Figure 2e–g presenting the scalar
 109 component of the electrical near-field magnitudes $\text{Re}\{|E_z|$,
 110 $\text{Re}\{|E_x|$, and $\text{Re}\{|E_y|$, respectively.
 111

The analysis of a nanoscale wireless communication system
 112 that transmits and receives optical plasmon polaritons with a
 113 cross section of $20\text{ nm} \times 20\text{ nm}$ to distance of $12\text{ }\mu\text{m}$ is shown
 114 in Figure 3. The system is comprised of a plasmonic
 115 nanotransceiver (Figure 1b) and a nanoplasmonic receiver,
 116 separated by a distance of $12\text{ }\mu\text{m}$. Figure 3a presents 3D AFM
 117 topography mapping of the nanosystem, where the transceiver
 118 is fabricated at the right-hand side and the receiver is located at
 119 the left side. To image the near-field structure of long-range
 120 plasmon polaritons transfer in real space, we use a combination
 121 of KPFM and SNOM. KPFM enables near-field mapping of
 122 plasmonic devices with a very high resolution. However, it has
 123 limited efficiency in characterizing dielectric devices, mainly
 124 since the work function of dielectric materials barely can be
 125 modified by optical illumination.^{3,4,6,12} Figure 3b shows KPFM
 126 analysis of the nanosystem, illuminated by a He–Ne laser ($\lambda =$
 127 633 m), linearly polarized in parallel to the dipole orientation
 128 and focused to diameter of 700 nm. Channel SPPs are observed
 129 at the nanotransceiver channel waveguide, propagate for 130

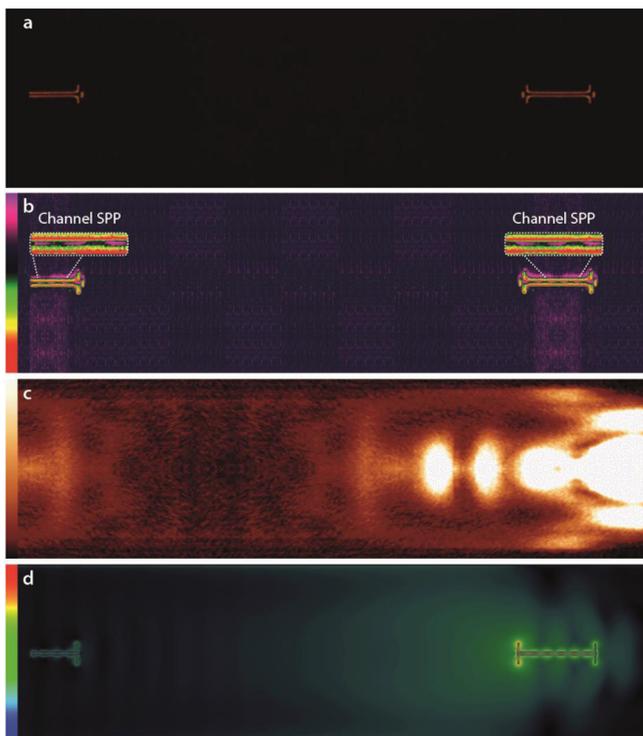


Figure 3. Characterization of the proposed wireless link, demonstrating efficient long-range propagation of tightly confined optical plasmon polaritons. (a) 3D AFM image of the fabricated wireless link system. (b) KPFM under optical illumination analysis of the wireless link system, showing SPPs at the transmission (right) and reception (left) sides; inset—zoom in to the respective waveguide channel. KPFM signal scale bar: ± 4.7 V. (c) SNOM analysis of the wireless link system, showing SPPs at the transmission (right) and reception (left) sides, as well as the coupling to surface waves that enable the long-range propagation. (d) Numerically calculated near-field image, showing the complete optical wireless transfer link. Scale bar: 750 nm.

distance of $L_{\text{WG}} = 1 \mu\text{m}$, followed by strong “hot spot” at the Tx 131
Plantenna that converts them to surface waves. A zoom in to 132
the channel region is presented in the inset, clearly showing the 133
periodic structure of the excited SPPs. Remarkably, pronounced 134
SPP excitation is observed at the distanced receiver, which is 135
not illuminated by the laser. Highlighted in the left inset, the 136
channel SPPs at the receiver waveguide are excited by efficient 137
coupling of surface waves to SPP by the receiver Plantenna. 138
The surface waves on the SiO_2 surface are imaged in the near 139
field via SNOM, as shown in Figure 3c. Naturally, SNOM 140
provides lower resolution images compared to KPFM,³ 141
however, its direct optical imaging mechanism enables mapping 142
of the surface photons that propagate on the dielectric medium, 143
unlike KPFM. Note that the SNOM image exhibits high 144
intensity at the physical locations of the transmitter and 145
receiver, originated by plasmon excitation. Hence, we state that 146
the combination of KPFM and SNOM provides a comple- 147
mentary, complete real-space nanoimaging approach for the 148
characterization of nanoscale wireless communication systems, 149
which facilitates high-resolution nanoimaging of both plasmons 150
and optical surface waves. Numerical calculation results of the 151
nanosystem, presenting the electric near-field magnitude $|\mathbf{E}|$, are 152
shown in Figure 3d. Both channel SPPs as well as the surface 153
waves in the dielectric substrate are clearly captured, providing 154
additional confirmation to our approach. 155

To unambiguously demonstrate the excellent efficiency of 156
our Plantenna based nanosystem, we compare its performances 157
to direct channel waveguiding link^{6–8,13} and to wireless link 158
based on dipole nanoantennas.^{14–17} For the wireless link 159
configurations (Figure 4a–d) the distance between the 160 f4
transceiver and receiver is $10 \mu\text{m}$, and for the direct link 161
(e.g., Figure 4e–f) the waveguide length is $3 \mu\text{m}$, limited by 162
fabrication constraints. Figure 4a shows AFM image of our 163
proposed Plantenna based nanosystem, as the corresponding 164
KPFM mapping is shown in Figure 4b with a voltage scale bar 165
of ± 4.7 V. Pronounced plasmon excitation is probed at the 166
receiver, evidenced by the modal structure of the field inside 167

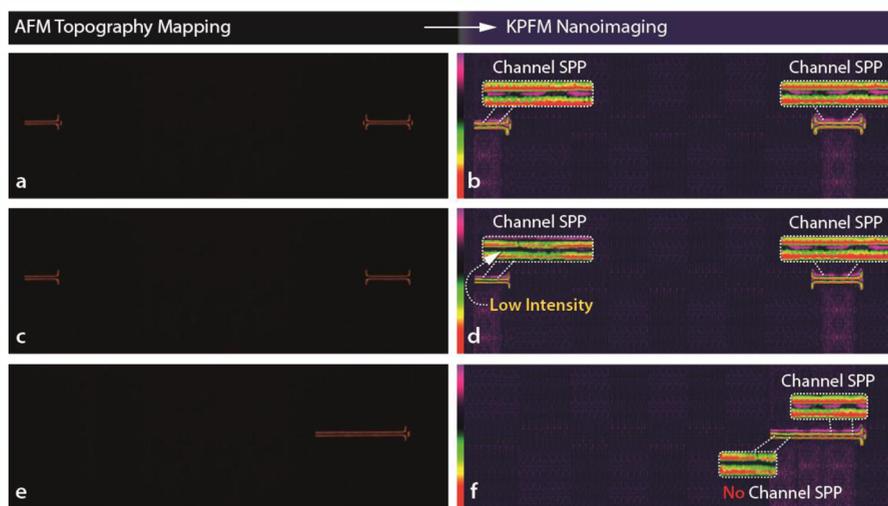


Figure 4. Comparison between nanophotonic links. (a) 3D AFM image of the Plantenna-based wireless link system. (b) KPFM under optical illumination analysis of the Plantenna-based wireless link system; KPFM signal scale bar: ± 4.7 V. (c) 3D AFM image of the dipole nanoantenna-based wireless link system. (d) KPFM under optical illumination analysis of the dipole nanoantenna-based wireless link system; KPFM signal scale bar: ± 0.5 V. (e) 3D AFM image of channel SPP waveguide with an identical cross section to the waveguides in a–d and $3 \mu\text{m}$ length. (f) KPFM under optical illumination analysis of the channel SPP waveguide link; KPFM signal scale bar: ± 4.7 V; inset (b, d, f): zoom in to the respective waveguide channel. Scale bar: $1 \mu\text{m}$.

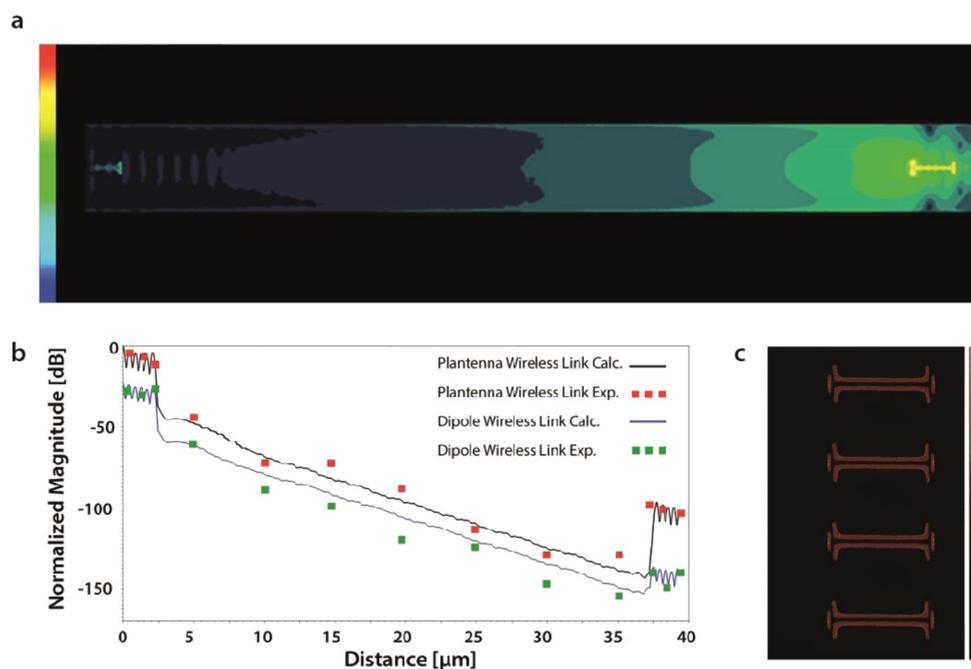


Figure 5. Performance analysis for ultralong propagation distances. (a) Numerically calculated near field image, showing the complete Plantenna-based wireless transfer link with distance of $35\ \mu\text{m}$ between the transmitting and receiving sides. (b) Performance comparison between Plantenna and dipole based wireless links for $35\ \mu\text{m}$, showing a 30 dB better performance of the Plantenna configuration. (c) 3D AFM image of a three-element Plantenna based transceiver phased array. Scale bar: 200 nm.

168 the channel which is highlighted in the inset. Figure 4c shows
 169 the AFM topography of a wireless link based on dipole
 170 nanoantennas, which was recently proposed as an approach for
 171 plasmonic energy transfer;¹⁸ the corresponding KPFM image is
 172 presented in Figure 4d with a voltage scale bar of $\pm 0.5\ \text{V}$.

173 We observe plasmon excitation at the transceiver; however,
 174 significantly less noticeable intensity is measured at the receiver
 175 waveguide (see inset) compared with the Plantenna based
 176 architecture. Figure 4e shows a 3D AFM image of a Plantenna
 177 integrated with a similar waveguide of $3\ \mu\text{m}$ length,
 178 implementing a direct nanoplasmonic link. Unlike the wireless
 179 links, the waveguide exhibits much higher propagation loss
 180 since it directly propagates tightly confined plasmons that
 181 interacts with the metals in their entire guided route.¹⁸ A
 182 KPFM map of the direct link is shown in Figure 4f (scale bar
 183 $\pm 4.7\ \text{V}$), where the zoom in to the different channel regions is
 184 presented in the insets. As seen in the right inset, channel SPPs
 185 are excited by the Plantenna and propagate through the
 186 waveguide. However, the huge propagation loss makes the
 187 waveguide SPPs decay significantly and being practically
 188 unobservable after propagating for only $2.5\ \mu\text{m}$, as seen in
 189 the left inset of Figure 4f. This reconfirms the critical, huge
 190 losses exhibited in gap plasmon waveguides with nanoscale
 191 channels, which hamper their real life applicability. Figure 5a
 192 shows the calculated electric near field for a Plantenna based
 193 communication nanosystem with a $35\ \mu\text{m}$ distance between the
 194 transmitter and receiver. Even for this high separation, plasmon
 195 excitation is clearly observed in the receiving device. Dipole
 196 nanoantenna is the most popular form of compact couplers to
 197 channel waveguide, enabling us to achieve 200 times higher
 198 efficiency compared with the case of directly illuminating a base
 199 waveguide.^{3,4,6,7,10,16} Other types of couplers may use nano-
 200 focusing approaches¹⁵ or more complex coupling devices like a
 201 Yagi nanoantenna.¹⁹ However, these devices are diffraction
 202 limited and cause only marginal improvement compared with

the dipole coupler. A comparison between Plantenna and 203
 dipole based systems with $35\ \mu\text{m}$ separation is shown in Figure 204
 5b, presenting the normalized near field along a line that 205
 connects the transmitter and receiver passing through the 206
 centers of both waveguides. The continuous black and blue 207
 charts represent the calculated results of a Plantenna and dipole 208
 wireless systems, respectively, as the discrete red and green 209
 squared dots are the corresponding experimental results. High 210
 field values are observed in both Tx and Rx ends, attributed to 211
 plasmon enhancement by the nanoantennas. As expected, the 212
 field is attenuated linearly when propagates through the SiO_2 213
 substrate. 214

A quantitative comparison between Plantenna and dipole 215
 systems is performed by comparing the SPP magnitude at both 216
 receiver waveguides, which serve as input for remote nano- 217
 plasmonic circuits^{4,12,20} or can be probed by photoelectric 218
 detectors. Since both systems are excited with identical sources 219
 and use similar plasmon waveguides, this approach is equivalent 220
 to calculating the ratio between the wave power of the SPP at 221
 the receiver waveguide and the laser source. As seen in Figure 222
 5b, the Plantenna based nanosystem outperforms the dipole 223
 configuration by more than 30 dB. Note that from nano- 224
 fabrication considerations we use Plantenna with identical 225
 dimensions through all of the experiments herein. However, the 226
 additional significant efficiency improvement can be achieved 227
 using a structural optimization of the different Plantennas as we 228
 recently reported.⁵ The Plantenna nanosystem can be used to 229
 wirelessly transfer optical nanoplasmonic information for 230
 macroscale distances using a phased array^{21–24} configuration 231
 as shown in Figure 5c. By fabricating an array of identical 232
 transceivers spaced by $\lambda_0/2$ (λ_0 is the free space wavelength), 233
 the emitted surface waves coupled from all of the transceivers 234
 can be coherently combined on the surface. Based on the well- 235
 known Friis principle,^{18,25,26} the phased array architecture 236
 enables propagation distances which are linearly scalable with 237

the number of transceivers, paving the way toward efficient wireless nanoplasmonic data and energy transfer for millimeter distances and beyond.

In conclusion, we designed, fabricated, and experimentally characterized a novel high-efficiency nanosystem, capable to wirelessly transfer deeply confined optical plasmon polaritons for chip-scale distances. Our system is architected for the efficient conversion of nanoscopic SPPs to propagating surface waves and to re-excite SPPs from these surface waves at significantly remote distances. We demonstrate the transmission of optical SPPs in channel waveguides with a cross section of only $20\text{ nm} \times 20\text{ nm}$, for distances which are 3 orders of magnitude larger than the plasmon wavelength. On the basis of the Plantenna, a new generation of high-performance nanoantennas with no RF equivalents, our nanosystem hugely outperforms both direct and wireless links based on dipole nanoantennas by more than 30 dB. For the first time, we use a unique combination of scanning probe microscopies to create complete real-space near-field mapping of a long-range nanoplasmonic wireless link at a high spatial resolution. This nanoimaging amalgamation provides valuable synergy needed for mapping both nanoscopic plasmon polaritons as well as macroscopically propagating surface waves. In the quest for reconciling the dimensional mismatch between diffraction-limited photonics and integrated electronics, our results enable new horizons for high integration densities of optical functionalities and interconnects. By using phased array configuration and utilizing degrees of freedom in polarization, frequency and code domains, inter- and intra-chip communications based on ultrafast nanoscale light waves as information carriers can now achieve record performances in terms of speed distance and size. The presented approach of hybridizing Plantenna and channel plasmon waveguides as nanotransceivers is immediately applicable for exploring long-range interaction between single and multiple quantum emitters, while our nanoimaging methodology enables enhanced understanding of exciting near-field phenomena at the nanoscale.

Methods. AFM and KPFM Measurements. All measurements were performed at room temperature and free ambient conditions (no vacuum), using a Dimension Icon AFM system with a NanoScope V controller (Bruker). For both AFM and KPFM measurements, we used NanoWorld probes SSS-NCH (SuperSharpSilicon, Noncontact/Tapping mode, High resonance frequency), with a typical diameter of 2 nm, resonance frequency of 320 kHz, and spring constant of 42 N/m. Typically, voltages of 2 V, AC capacitance frequencies of 880 MHz, lift heights of 30–50 nm, and line rates of 0.1 kHz were employed. To map the CPD of the sample, we apply both AC voltage (VAC) and DC voltage (VDC) to the AFM tip. VAC generates oscillating electrical forces between the AFM tip and sample surface, and VDC nullifies the oscillating electrical forces that originated from CPD between tip and sample surface.

Optical Near-Field Measurements. The optical characterization of the plasmonic structures was performed by a MultiView 2000 scanning probe microscope/NSOM system (Nanonics Imaging Ltd.). The SPM head was placed on the stage of an Olympus dual microscope while remaining the optical axis free from above and below. Such a configuration allowed us to bring the cantilevered NSOM tip to the desired position on the sample under an upper objective of 50X. The sample was illuminated with a Liconix diode laser of 785 laser

CW light from the bottom and focused on a sample with a $50\times$ objective. We used the bottom piezo scanner of the scanning head to place the desired structures of the sample very accurately relatively to the incoming light of the laser from below. The scan was performed with upper piezo scanner allowing moving only the NSOM tip while the sample remains still. The collection of near field light distribution on the surface was performed in tapping mode with a 200 nm aperture NSOM tips based tuning fork produced by Super Tips (Nanonics Imaging Ltd.). The signal was transmitted through a multimode optical fiber onto an APD. The AFM and NSOM images were collected simultaneously during the scan, allowing to monitor the topography of the desired structure and to correlate it with the near-field optical signal that comes from any particular feature.

Numerical Simulations. The numerical results are obtained by using the software package ANSYS HFSS V15, the industry-standard simulation tool for 3D full-wave electromagnetic field simulation. HFSS solves Maxwell's equations via the finite element method (FEM) using an adaptive mesh refinement process for tailored accuracy requirements. The field's solutions are calculated with the metallic (Ag) plasmonic structures being deposited on a homogeneous SiO_2 substrate. The nanoantenna is illuminated by optical sources at 474 THz (wavelength of 633 nm), which are modeled as focused Gaussian beams with 1 μm characteristic diameter. The electric field is polarized in parallel with the dipole direction, as the wave vector \mathbf{K} is perpendicular. A selectively dense meshing is assigned in the metallic and waveguiding regions, with a maximum cell size of 1 nm and 750 000 FEM tetrahedral cells. To provide maximum accuracy, the model is terminated as following: the interface with free space is bounded by perfectly matched layer (PML) absorbing boundary conditions (ABC), while the metallic and SiO_2 termination are done via layered impedance (LI) ABC. The minimum number of adaptive meshing iterations was set to 12, with a convergence condition of 1% maximum energy variance between adjacent iterations.

Fabrication. SiO_2/Si sample was spin-coated with poly(methyl methacrylate) (PMMA 950 A2) electron-beam resist providing thickness of 100 nm. The samples coated with PMMA were subsequently baked for 120 s on a hot plate at 180C. The desired pattern was exposed in the PMMA layer using a CRESTEC CABLE-9000C high-resolution electron-beam lithography system using different doses to control line and gap width. Then the samples were developed for 90 s using methyl isobutyl ketone (MIBK) and rinsed with IPA. The samples were subsequently exposed to Ar plasma to etch 10 nm in order to remove leftovers from the pattern, sputtered using BESTEC 2" DC magnetron to deposit 2 nm Cr and 18 nm Au, and then immersed in 180 KHz ultrasonic bath with NMP for 3 h for resist liftoff.

AUTHOR INFORMATION

Corresponding Author

*E-mail: moshik.cohen80@gmail.com.

ORCID

Moshik Cohen: [0000-0002-7519-9742](https://orcid.org/0000-0002-7519-9742)

Author Contributions

M.C. carried out the theoretical design and analysis, designed the studies, performed the experiments, and wrote the manuscript. Z.Z. and R.S. participated in writing the manuscript and designing the study.

362 **Notes**

363 The authors declare no competing financial interest.

364 ■ **ACKNOWLEDGMENTS**

365 M.C. acknowledges Olga Girshevitz Yafit from Bar-Ilan
366 Institute for Nanotechnology & Advanced Materials (BINA),
367 for the support in fabricating the reported structures.

368 ■ **REFERENCES**

- 369 (1) Schuller, J. A.; Barnard, E. S.; Cai, W.; Jun, Y. C.; White, J. S.;
370 Brongersma, M. L. *Nat. Mater.* **2010**, *9* (3), 193–204.
- 371 (2) Brongersma, M. L.; Shalae, V. M. *Science* **2010**, *328* (5977),
372 440–441.
- 373 (3) Cohen, M.; Shavit, R.; Zalevsky, Z. *Sci. Rep.* **2015**, *4*, 04096.
- 374 (4) Cohen, M.; Zalevsky, Z.; Shavit, R. *Nanoscale* **2013**, *5* (12),
375 5442–5449.
- 376 (5) Cohen, M.; Shavit, R.; Zalevsky, Z. *Sci. Rep.* **2015**, *5*, 17562.
- 377 (6) Cohen, M.; Shavit, R.; Zalevsky, Z. In *Planar Waveguides and*
378 *other Confined Geometries*; Marowsky, G., Ed.; Springer Series in
379 Optical Sciences; Springer: New York, 2015; pp 45–66.
- 380 (7) Dionne, J. A.; Sweatlock, L. A.; Atwater, H. A.; Polman, A. *Phys.*
381 *Rev. B: Condens. Matter Mater. Phys.* **2006**, *73* (3), 035407.
- 382 (8) Bozhevolnyi, S. I.; Volkov, V. S.; Devaux, E.; Laluet, J.-Y.;
383 Ebbesen, T. W. *Nature* **2006**, *440* (7083), 508–511.
- 384 (9) Dregely, D.; Lindfors, K.; Lippitz, M.; Engheta, N.; Totzeck, M.;
385 Giessen, H. *Nat. Commun.* **2014**, *5*, 4354.
- 386 (10) Cohen, M.; Abulafia, Y.; Shavit, R.; Zalevsky, Z. *ACS Nano*
387 **2017**, *11*, 3274.
- 388 (11) Carmeli, I.; Cohen, M.; Heifler, O.; Lilach, Y.; Zalevsky, Z.;
389 Mujica, V.; Richter, S. *Nat. Commun.* **2015**, *6*, 7334.
- 390 (12) Cohen, M.; Shavit, R.; Zalevsky, Z.; Abulafia, Y. In *CLEO: 2014*;
391 OSA Technical Digest (online); Optical Society of America, 2014; p
392 JTu4A.135.
- 393 (13) Lu, H.; Liu, X.; Wang, G.; Mao, D. *Nanotechnology* **2012**, *23*
394 (44), 444003.
- 395 (14) Wang, S.; Zhan, Q. *Sci. Rep.* **2016**, *6*, 29626.
- 396 (15) Choo, H.; Kim, M.-K.; Staffaroni, M.; Seok, T. J.; Bokor, J.;
397 Cabrini, S.; Schuck, P. J.; Wu, M. C.; Yablonovitch, E. *Nat. Photonics*
398 **2012**, *6* (12), 838–844.
- 399 (16) Wen, J.; Romanov, S.; Peschel, U. *Opt. Express* **2009**, *17* (8),
400 5925–5932.
- 401 (17) Andryeuskii, A.; Zenin, V. A.; Malureanu, R.; Volkov, V. S.;
402 Bozhevolnyi, S. I.; Lavrinenko, A. V. *Nano Lett.* **2014**, *14* (7), 3925–
403 3929.
- 404 (18) Alù, A.; Engheta, N. *Phys. Rev. Lett.* **2010**, *104* (21), 213902.
- 405 (19) Solís, D. M.; Taboada, J. M.; Obelleiro, F.; Landesa, L. *Opt.*
406 *Express* **2013**, *21* (2), 2369–2377.
- 407 (20) Dionne, J. A.; Sweatlock, L. A.; Sheldon, M. T.; Alivisatos, A. P.;
408 Atwater, H. A. *IEEE J. Sel. Top. Quantum Electron.* **2010**, *16* (1), 295–
409 306.
- 410 (21) Abediasl, H.; Hashemi, H. *Opt. Express* **2015**, *23* (5), 6509–
411 6519.
- 412 (22) Cohen, E.; Ruberto, M.; Cohen, M.; Degani, O.; Ravid, S.;
413 Ritter, D. *IEEE Trans. Microwave Theory Tech.* **2013**, *61* (3), 1359–
414 1375.
- 415 (23) Sun, J.; Timurdogan, E.; Yaacobi, A.; Hosseini, E. S.; Watts, M.
416 R. *Nature* **2013**, *493* (7431), 195–199.
- 417 (24) Maguid, E.; Yulevich, I.; Veksler, D.; Kleiner, V.; Brongersma,
418 M. L.; Hasman, E. *Science* **2016**, *352*, 1202.
- 419 (25) Choi, H. S.; Kang, S. Y.; Cho, S. J.; Oh, I.-Y.; Shin, M.; Park, H.;
420 Jang, C.; Min, B.-C.; Kim, S.-I.; Park, S.-Y.; Park, C. S. *Sci. Rep.* **2015**, *4*,
421 5486.
- 422 (26) Yang, Y.; Li, Q.; Qiu, M. *Sci. Rep.* **2016**, *6*, 19490.